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LETTER TO THE EDITOR

Boson realisation of Virasoro and Kac–Moody algebras and their indecomposable representations

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Abstract. In this letter we give the boson realisation of Virasoro and Kac–Moody algebras without central terms and construct their indecomposable representations from some representations of the Heisenberg–Weyl algebra. The boson realisation of the Kac–Moody algebra $\hat{\mathfrak{su}}(2)$ associated to the Lie algebra $\mathfrak{su}(2)$ is discussed in detail.

Virasoro and Kac–Moody algebras have appeared in many areas of physics [1]. In superstring theory the Virasoro algebra plays an important role [2]. The Kac–Moody symmetries are found in (1+1)-dimensional classical field theory, Yang–Mills theory and so on. More recently, the new Virasoro and Kac–Moody symmetries in the non-linear σ model [3] and the solution space of the Ernst equation [4] have been studied by Hou and Li.

Although the representation theory of Virasoro and Kac–Moody algebras have been studied by many authors (e.g. [5]), we will pay attention to their physical indecomposable representations. By our method used to study the Lie algebras [6] and Lie superalgebras [7], we discuss Virasoro and Kac–Moody algebras in this letter.

On the quotient space $\Omega = \bar{\Omega}/L$:

$$\left\{ f(k_i, S_i) = \prod_{i=1}^N (b_i^{+k_i} b_i^{S_i}) \text{ mod } L \mid S_i, k_i \in \mathbb{N}, i = 0, 1, 2, \dots, N \right\} \quad (1)$$

of the universal enveloping algebra $\bar{\Omega}$ of the Heisenberg–Weyl algebra $H : \{b_i^+, b_i, e\}$, where L is a left ideal generated by the element $e - 1$, an indecomposable representation of H is obtained in [6] as

$$\begin{aligned} \rho(b_i^+)f(k_i, S_i) &= f(k_i + \delta_{ii}, S_i) \\ \rho(b_i)f(k_i, S_i) &= f(k_i, S_i + \delta_{ii}) + k_i f(k_i - \delta_{ii}, S_i) \\ \rho(e)f(k_i, S_i) &= f(k_i, S_i). \end{aligned} \quad (2)$$

From the above representation, we will construct certain types of representation of Virasoro and Kac–Moody algebras.

Let g be a Lie algebra with generators $\{T^a, a = 1, 2, \dots, M\}$ that satisfy the commutation relations

$$[T^a, T^b] = \sum_{c=1}^M \mathcal{F}_c^{ab} T^c \quad (3)$$

where complex numbers \mathcal{G}_c^{ab} are the structure constants of Lie algebra g . For each given N -dimensional representation $P = [P_{ij}]$, ($i = 1, 2, \dots, N$) of Lie algebra g , we define

$$T_l^a = b_0^{+l} [b_1^+, b_2^+, \dots, b_N^+] [P_{ij}] \begin{bmatrix} b_1 \\ b_2 \\ \vdots \\ b_N \end{bmatrix} \tag{4}$$

where because of the existence of the left inverse of b_0^+ , l can be taken as a integer, i.e. $l \in \mathbb{Z}$.

These T_l^a form an untwisted affine Kac-Moody algebra \hat{g} without a central term. In fact, considering the explicit expression of T_l^a

$$T_l^a = b_0^{+l} \sum_{m,n} P_{mn}(T^a) b_m^+ b_n \tag{5}$$

we easily prove

$$[T_l^a, T_h^b] = \sum_{c=1}^M \mathcal{G}_c^{ab} T_{l+h}^c \tag{6}$$

Defining the operators

$$L_l = -b_0^{l+1} b_0 \tag{7}$$

we obtain a Virasoro algebra $\hat{v} : \{L_l | l \in \mathbb{Z}\}$ that satisfies

$$[L_l, L_h] = [l-h] L_{l+h} \tag{8}$$

$$[L_l, T_h^a] = -h T_{l+h}^a \tag{9}$$

Then (4) and (7) give the boson realisation of the Kac-Moody and Virasoro algebras.

According to (2), (4) and (7), extending the quotient space Ω to $k_0 \in \mathbb{Z}$, we obtain the representations

$$\begin{aligned} \Gamma(T_l^a) f(k_i, S_i) &= \sum_{m,n} P_{mn}(T^a) f(k_i + l\delta_{i0} + \delta_{mi}, S_i + \delta_{in}) \\ &+ \sum_{mn} P_{mn}(T^a) f(k_i - \delta_{in} + l\delta_{i0} + \delta_{im}, S_i) \end{aligned} \tag{10}$$

$$\Gamma(L_l) f(k_i, S_i) = f(k_i + \delta_{i0} l, \delta_i + \delta_{i0}) + k_0 f(k_i + l\delta_{i0}, S_i) \tag{11}$$

of the Kac-Moody and Virasoro algebras. By the same analysis as that in [6, 7], we see that the above representations are indecomposable.

On the quotient space $V = \Omega / L_+$

$$\{F(k_i) = F(k_i, 0) \text{ mod } L_+\}$$

where the left ideal L_+ is generated by elements $b_i - \Lambda_i$ ($i = 1, 2, \dots, N$; $\Lambda_i \in \mathbb{C}$) the representations (10) and (11) induce the new indecomposable representation

$$\begin{aligned} \Gamma(T_l^a) F(k_i) &= \sum_{m,n} P_{mn}(T^a) \Lambda_n F(k_i + l\delta_{i0} + \delta_{mi}) + \sum_{mn} k_n P_{mn}(T^a) F(k_i + l\delta_{i0} - \delta_{in} + \delta_{im}) \\ \Gamma(L_l) \Gamma(k_i) &= \Lambda F(k_i + \delta_{i0} l) + k_0 F(k_i + l\delta_{i0}) \end{aligned} \tag{12}$$

for $\Lambda_i \neq 0$.

For the case with $\Lambda_i = 0$ ($i = 0, 1, \dots, N$), the representations

$$\begin{aligned} \Gamma(T_i^a)F(k_i) &= \sum_{m,n} k_n P_{mn}(T^a)F(k_i + l\delta_{i0} - \delta_{mi} + \delta_{ni}) \\ \Gamma(L_i)F(k_i) &= \Lambda_0 F(k_i + l\delta_{i0}) \end{aligned} \tag{13}$$

given by (12) are completely reducible. This is due to the invariance of the sum $\sum_{i=1}^N k_i$ under the action of the representation (13).

Following the above general discussion, we study the Kac-Moody algebra $\widehat{su}(2)$ associated with the Lie algebra $su(2)$. According to (4), we have the boson realisation of $\widehat{su}(2)$

$$\begin{aligned} J_i^+ &= b_0^{+l} b_1^+ b_2 & J_i^- &= b_0^{+l} b_2^+ b_1 \\ J_i^3 &= \frac{1}{2} b_0^{+l} [b_1^+ b_1 - b_2^+ b_2] \end{aligned} \tag{14}$$

corresponding to the Pauli representation of $su(2)$. The boson realisation of an indecomposable representation of $\widehat{su}(2)$ is obtained as

$$\begin{aligned} \Gamma(J_i^+)F(k_0, k_1, k_2) &= \Lambda_2 F(k_0 + l, k_1 + 1, k_2) + k_2 F(k_0 + l, k_1 + 1, k_2 - 1) \\ \Gamma(J_i^-)F(k_0, k_1, k_2) &= \Lambda_1 F(k_0 + l, k_1, k_2 + 1) + k_1 F(k_0 + l, k_1 - 1, k_2 + 1) \\ \Gamma(J_i^3)F(k_0, k_1, k_2) &= \frac{1}{2} \Lambda_1 F(k_0 + l, k_1 + 1, k_2) - \frac{1}{2} \Lambda_2 F(k_0 + l, k_1, k_2 + 1) \\ &\quad + \frac{1}{2} (k_1 - k_2) F(k_0 + l, k_1, k_2). \end{aligned} \tag{15}$$

When $\Lambda_0 = \Lambda_1 = \Lambda_2 = 0$, define a new basis for the space V

$$\tilde{F}[M, n, k] = F(k, n, M - n) \quad k \in \mathbb{Z}; n, M \in \mathbb{N}.$$

On this space, we obtain a representation of $\widehat{su}(2)$

$$\begin{aligned} \Gamma(J_i^+) \tilde{F}[M, n, k] &= (M - n) \tilde{F}[M, n + 1, k + l] \\ \Gamma(J_i^-) \tilde{F}[M, n, k] &= n \tilde{F}[M, n - 1, k + l] \\ \Gamma(J_i^3) \tilde{F}[M, n, k] &= (2n - M) \tilde{F}[M, n, k + l]. \end{aligned} \tag{16}$$

Then, for fixed $M \in \mathbb{N}$, $\{\tilde{F}(M, n, k) | n \in \mathbb{N}, k \in \mathbb{Z}\}$ forms an invariant subspace $V^{(M)}$ of V , and V can be written as the direct sum of spaces, i.e.

$$V = V^{(0)} \oplus V^{(1)} \oplus V^{(2)} \oplus \dots \oplus V^{(N)} \oplus.$$

Thus, for the case with $\Lambda_0 = \Lambda_1 = \Lambda_2 = 0$, the representation (16) is completely reducible. On each space $V^{(M)}$, (16) gives an irreducible representation of $\widehat{su}(2)$.

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